

Proton Decay Matrix Elements from Lattice QCD

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for RBC/UKQCD collaboration

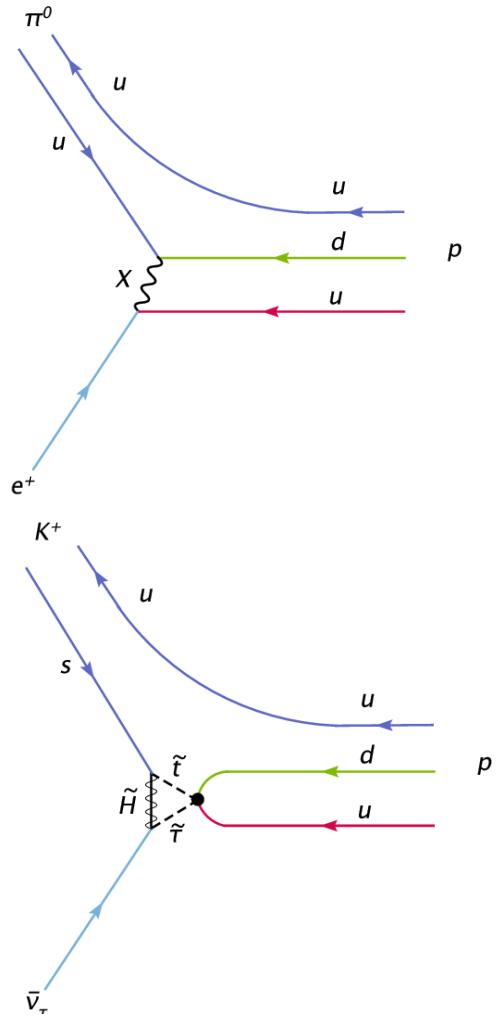
Proton decay

- Baryon asymmetric process
 - $p \rightarrow e^+ + \pi^0$, etc
 - Baryogenesis
 - GUTs scaling physics
(in the SM, $\tau(t \rightarrow e^+ \nu_\mu \nu_\tau) \sim 10^{150}$ yr by anomalies)
 - Mediated by colored and flavored heavy particle
- Low energy matrix element
 - $p \rightarrow [\text{PS} + \text{lep}]$ lifetime, Br, ...
needs QCD matrix element
 \Rightarrow model prediction

GUTs

- Baryon number breakings
 - Mediated by X boson
minimal SU(5), dim 6 Op
predicts $\tau(p \rightarrow \pi^0 e^+) \sim 10^{30} \text{ yr}$
 \Rightarrow ruled out by recent SK results
 - $O^i(\mu) = (qq)_\Gamma (ql)_{\Gamma'}$

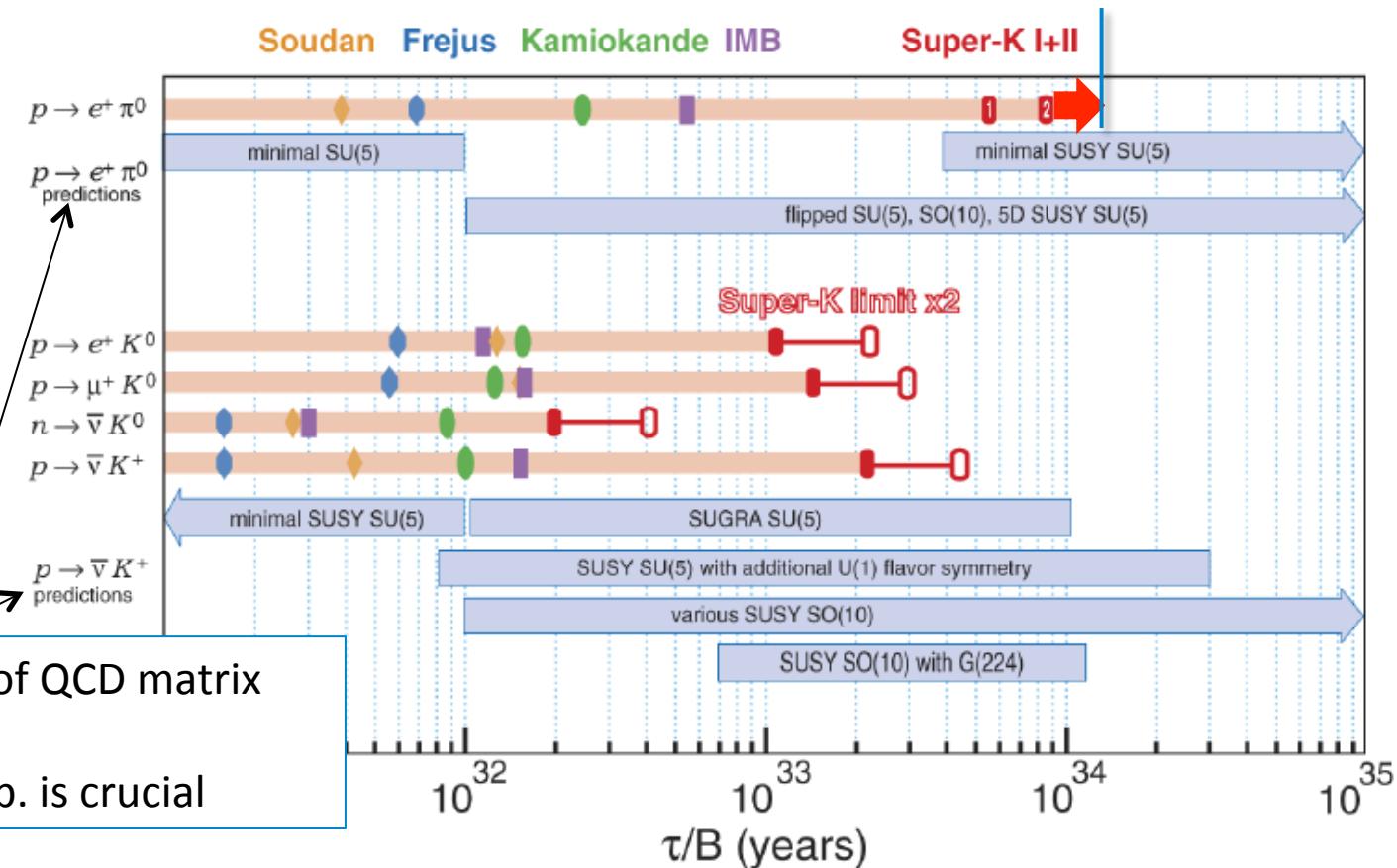
- Mediated by colored sparticle, Higgsino (or wino)
SUSY SU(5), dim 5 Op
increase lifetime, dominate
- $p \rightarrow K^+ + \bar{\nu}$ mode



Proton lifetime limit

- Recent Super-Kamiokande results:

$$\tau_{\text{exp}}(p \rightarrow \pi^0 e^+) > 1.01 \times 10^{34} \text{ yr} (\text{SK, 2009})$$



Ed Kearns (2007)

Matrix element

$$C_i(\mu)O_i(\mu)/M_{\text{GUT}}^2$$

- Expansion to dim 6 coupling

$$\mathcal{L}_{\text{GUT}} = \mathcal{L}_{\text{SM}} + \sum_i C_i(\mu) O_i(\mu) / M_{\text{GUT}}^2$$

C_i depends on type of GUTs model

$$\langle \pi^0 e^+ | p \rangle_{\text{GUT}} = \sum_{i=\Gamma, \Gamma'} C_i \langle \pi^0 | \mathcal{O}_{udu}^{\Gamma \Gamma'} | p \rangle \bar{v}_{e^+}^c$$

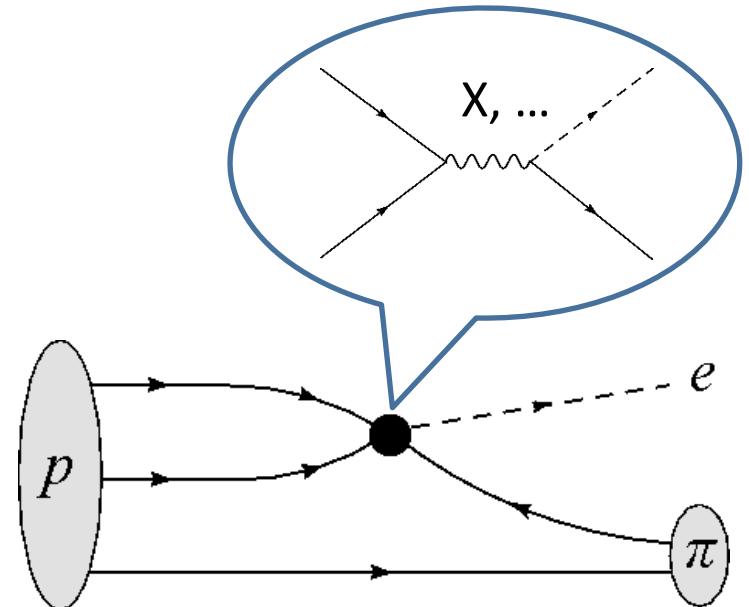
$\mathcal{O}_{abc}^{\Gamma \Gamma'} = (q_a q_b)_\Gamma q_c \Gamma' : 3 \text{ quark operator}$

- Decay width

$$\Gamma_{p \rightarrow \pi^0 e^+} = \frac{m_p}{32\pi^2} \left[1 - \left(\frac{m_e}{m_p} \right)^2 \right]^2 \left| \sum_i C_i W_0^i(p \rightarrow \pi^0) \right|^2$$

W_0^i : determine from QCD matrix element

⇒ lattice calculation is needed

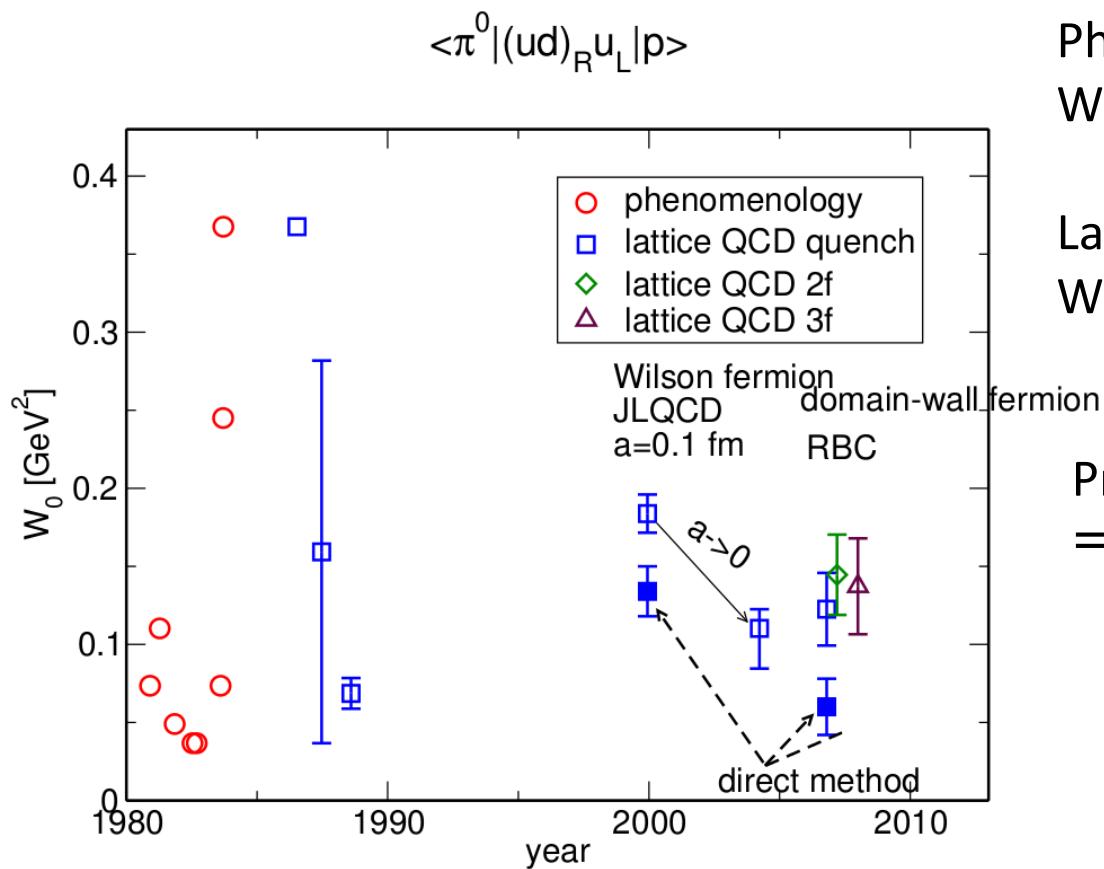


$$\langle PS; \vec{p} | \mathcal{O}_{abc}^{\Gamma \Gamma'} | N; \vec{k}, s \rangle \bar{v}_{e^+}^c$$

Lattice QCD

- The first principles of QCD

Determination of $[p \rightarrow \text{PS}]$ matrix element



Phenomenological estimate:
 $W_0 (p \rightarrow \pi) \simeq 0.1 -- 0.4 \text{ GeV}^2$

Lattice results:
 $W_0 (p \rightarrow \pi) \simeq 0.1 -- 0.2 \text{ GeV}^2$

Proton lifetime $\sim O(m^3/W_0^2)$
 \Rightarrow factors of $O(10)$ uncertainties

RBC/UKQCD efforts

- RBC(2007) Y. Aoki et al.
PRD75 014507, hep-lat/0607002
 - Direct/indirect method in Quench DW
 - Dynamical effect (indirect method)
 - Non-perturbative renormalization (NPR)
- RBC/UKQCD(2008) Y. Aoki et al.
PRD78 054505 arXiv:0806.1031
 - Indirect method in $N_f=3$ full QCD with dynamical DW
 - NPR
- RBC/UKQCD(2011) Y. Aoki et al. for RBC/UKQCD in prep.
 - Direct method in $N_f=3$ full QCD with dynamical DW
 - 2.4 fm^3 and pion mass: $315 -- 615 \text{ MeV}$
 - NPR

W_0 from lattice QCD

- Indirect method JLQCD(2000), Y. Aoki (2007)

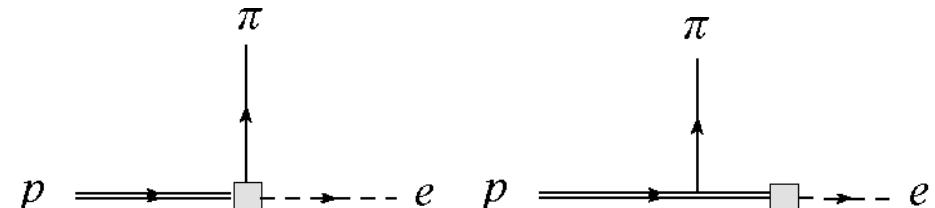
- Via soft-pion theorem,

$$\lim_{p \rightarrow 0} \langle \pi^k(p) | \mathcal{O}_{udu}^{\Gamma\Gamma'} | N \rangle = -\frac{i}{f} \langle 0 | [Q_5^k, \mathcal{O}_{udu}^{\Gamma\Gamma'}] | N \rangle$$

$$\langle 0 | (\bar{u}d)_L u | p \rangle = \alpha u_p, \quad \langle 0 | (\bar{u}d)_R u | p \rangle = \beta u_p$$

- Low-energy-constant is determined by 2pt function which easily and cheaply computes.
 - Through Chiral perturbation theory, matrix element

$$W_0^L(\pi^0 \rightarrow p) \simeq \frac{\alpha}{\sqrt{2}f} (1 + D + F)$$



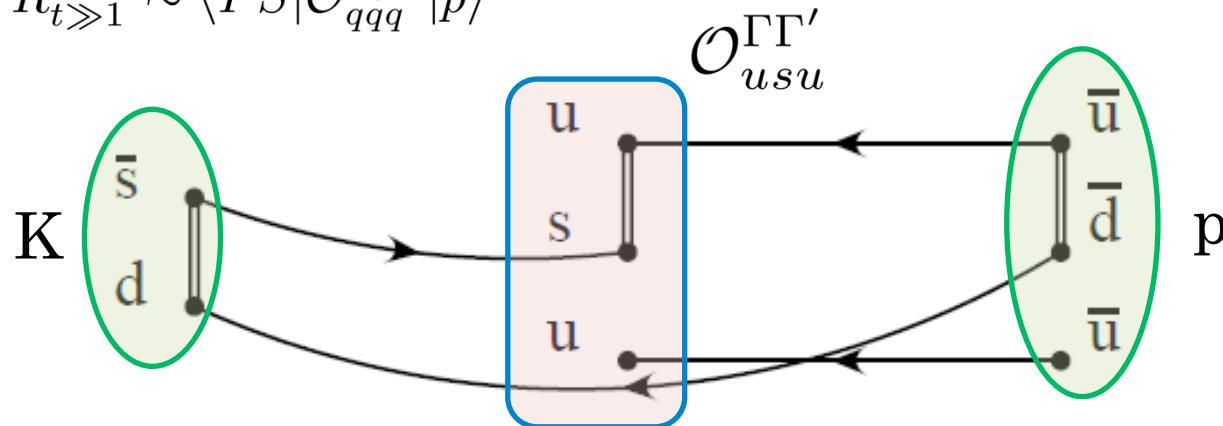
- Uncertainties remain in soft-pion theorem from large mass and momentum

W_0 from lattice QCD

- Direct method CP-PACS(2000), Y. Aoki (2007)
 - Computation of 3pt function including B violating operator
 - Ratio provides matrix element and W_0

$$R^{L/R} = \frac{\langle 0 | J_{PS}(\vec{p}, t_{PS}) \mathcal{O}_{qqq}^{L/R}(\vec{p}, t) J_p(\vec{0}, t_p) | 0 \rangle}{\langle 0 | J_p(\vec{0}, t_p) J_p(\vec{0}, 0) | 0 \rangle \langle 0 | J_{PS}(\vec{p}, t_{PS}) J_{PS}(\vec{p}, t) | 0 \rangle} \sqrt{Z_{PS} Z_p} V$$

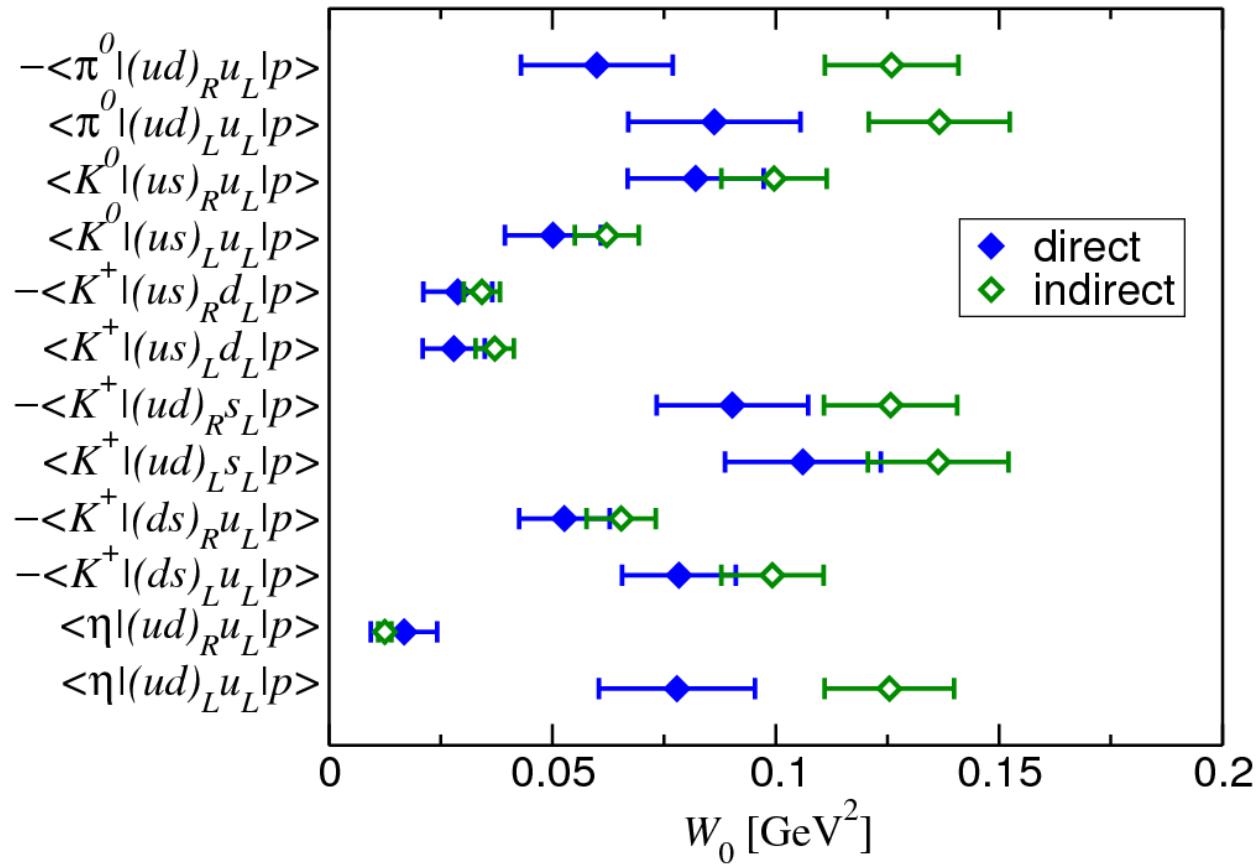
$$R_{t \gg 1}^{L/R} \sim \langle PS | \mathcal{O}_{qqq}^{L/R} | p \rangle$$



- No intermediate model including the above estimate
- Individually obtain W_0 which depends on final meson state.

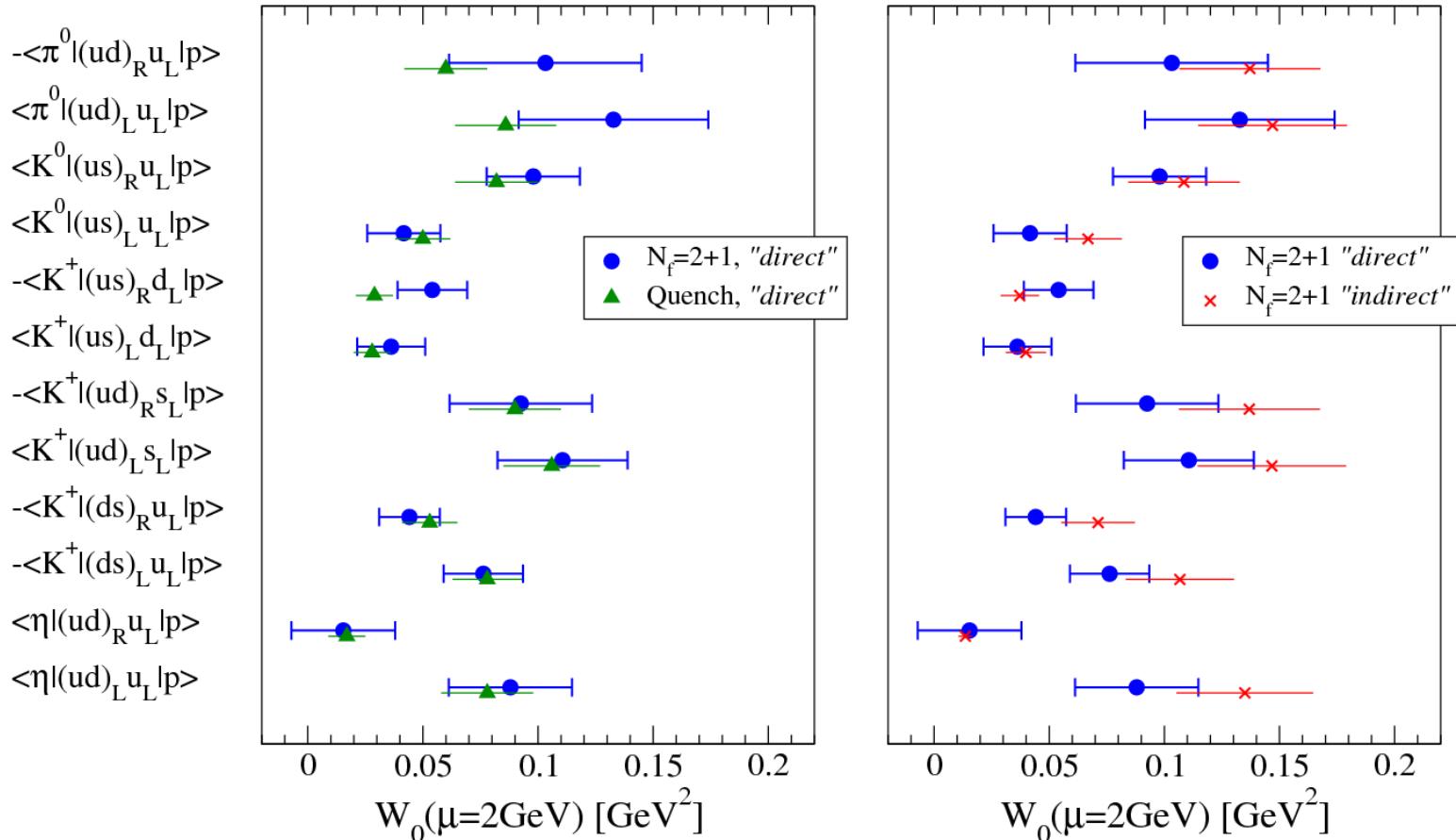
In quench QCD

Y. Aoki (2007)



In full QCD

RBC/UKQCD in prep.



- There is no significant discrepancy between each results.
- Statistical and systematic errors are still large.

Error budget

| | $W_0(\mu = 2\text{GeV})$ | GeV^2 | Extrapolation | $\mathcal{O}(a^2)$ | ΔZ | Δa^{-1} |
|---|--------------------------|----------------|---------------|--------------------|------------|-----------------|
| $-\langle \pi^0 (ud)_R u_L p \rangle$ | 0.103(23)(34) | | 0.033 | 0.005 | 0.008 | 0.004 |
| $\langle \pi^0 (ud)_L u_L p \rangle$ | 0.133(29)(28) | | 0.026 | 0.007 | 0.011 | 0.005 |
| $\langle K^0 (us)_R u_L p \rangle$ | 0.098(15)(12) | | 0.007 | 0.005 | 0.008 | 0.003 |
| $\langle K^0 (us)_L u_L p \rangle$ | 0.042(13)(8) | | 0.007 | 0.002 | 0.003 | 0.001 |
| $-\langle K^+ (us)_R d_L p \rangle$ | 0.054(11)(9) | | 0.008 | 0.003 | 0.004 | 0.002 |
| $\langle K^+ (us)_L d_L p \rangle$ | 0.036(12)(7) | | 0.007 | 0.002 | 0.003 | 0.001 |
| $-\langle K^+ (ud)_R s_L p \rangle$ | 0.093(24)(18) | | 0.016 | 0.005 | 0.008 | 0.003 |
| $\langle K^+ (ud)_L s_L p \rangle$ | 0.111(22)(16) | | 0.012 | 0.006 | 0.009 | 0.004 |
| $-\langle K^+ (ds)_R u_L p \rangle$ | 0.044(12)(5) | | 0.003 | 0.002 | 0.004 | 0.002 |
| $-\langle K^+ (ds)_L u_L p \rangle$ | 0.076(14)(9) | | 0.006 | 0.004 | 0.006 | 0.003 |
| $\langle \eta (ud)_R u_L p \rangle$ | 0.015(14)(17) | | 0.017 | 0.001 | 0.001 | 0.001 |
| $\langle \eta (ud)_L u_L p \rangle$ | 0.088(21)(16) | | 0.014 | 0.004 | 0.007 | 0.003 |

← uncertainty of extrapolation to physical pion mass is 30 % error, and statistical error is also 20 %.

Future plan

- Improvement of statistical errors Izubuchi and Shintani in prep.
 - New algorithm, LMA/AMA in which computational cost can be reduced by factor 1/10 -- 1/100.
- Improvement of systematic errors
 - Computation in more realistic lattice parameters
 - Reduction of chiral extrapolation uncertainty using small quark mass
⇒ ~10% accuracy

| Lattice size | Physical size | Lattice spacing | L_s | Gauge action | Pion mass |
|------------------|--------------------|-----------------|-------|--------------|----------------|
| $24^3 \times 64$ | 2.7 fm^3 | 0.114 fm | 16 | Iwasaki | 315 -- 615 MeV |
| $32^3 \times 64$ | 2.7 fm^3 | 0.087 fm | 16 | Iwasaki | 295 -- 397 MeV |
| $32^3 \times 64$ | 4.6 fm^3 | 0.135 fm | 32 | DSDR | 171 -- 241 MeV |
| $48^3 \times 96$ | 5.5 fm^3 | 0.115 fm | 8 | Iwasaki | 135 MeV |

Backup

non-perturbative renormalization

- * To avoid lattice perturbation theory: expansion poorly converges
- * RI/MOM scheme constructed: [YA et al., PRD75(07)014507] → NLO matching to \overline{MSbar}
- * Exact chiral symmetry → multiplicative renormalization
- * DWF has negligible chiral symmetry breaking
- * In general, mixing occurs

$$O_{RL} = (\bar{u^c} P_R d) \cdot P_L s$$

$$O_{LL} = (\bar{u^c} P_L d) \cdot P_L s$$

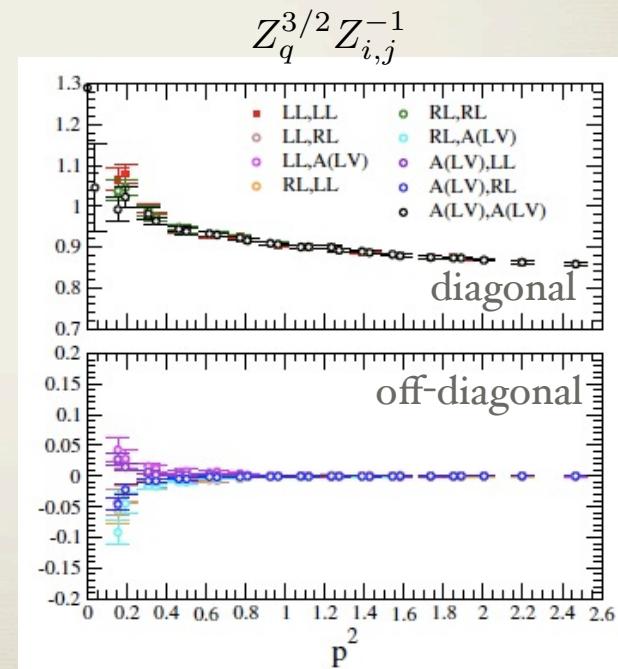
$$O_{A(LV)} = (\bar{u^c} \gamma_\mu \gamma_5 d) \cdot P_L \gamma_\mu s$$

- * Z : diagonal matrix: practically no mixing!

$$U^{\overline{MS} \leftarrow \text{latt}}(\mu = 2 \text{ GeV})_{LL} = 0.662 \pm 0.010$$

$$U^{\overline{MS} \leftarrow \text{latt}}(\mu = 2 \text{ GeV})_{RL} = 0.665 \pm 0.008,$$

+8% $\mathcal{O}(\alpha_s^2)$ error



[RBC/UKQCD YA et al. PRD78(08)04505]

P decay operator

- Weinberg rep.

$$\mathcal{O}_{abcd}^1 = (D_a^i, U_b^j)_R (q_c^{k\alpha}, l_d^\beta)_L \varepsilon^{ijk} \varepsilon^{\alpha\beta},$$

$$\mathcal{O}_{abcd}^2 = (q_a^{i\alpha}, q_b^{j\beta})_L (U_c^k, l_d)_R \varepsilon^{ijk} \varepsilon^{\alpha\beta},$$

$$\tilde{\mathcal{O}}_{abcd}^4 = (q_a^{i\alpha}, q_b^{j\beta})_L (q_c^{k\gamma}, l_d^\delta)_L \varepsilon^{ijk} \varepsilon^{\alpha\delta} \varepsilon^{\beta\gamma},$$

$$\mathcal{O}_{abcd}^5 = (D_a^i, U_b^j)_R (U_c^k, l_d)_R \varepsilon^{ijk}$$

a, b, c, d : generation, $\alpha, \beta, \gamma, \delta$: SU(2) indices

i,j,k: color indices

General form of three quark:

$$\mathcal{O}_{abc}^{\Gamma\Gamma'} = (q_a q_b)_\Gamma q_c \Gamma' = (q_a^{Ti} C P_\Gamma q_b^j) P_{\Gamma'} q_c^k \varepsilon^{ijk}$$

and matrix element

$$\langle PS; \vec{p} | \mathcal{O}_{abc}^{\Gamma\Gamma'} | N; \vec{k}, s \rangle \quad s: \text{nucleon spin}$$

P decay operator

- Parity invariance

$$\langle PS; \vec{p} | \mathcal{O}^{LL} | N; \vec{k}, s \rangle = \gamma_4 \langle PS; -\vec{p} | \mathcal{O}^{RR} | N; -\vec{k}, s \rangle, \quad \langle PS; \vec{p} | \mathcal{O}^{LR} | N; \vec{k}, s \rangle = \gamma_4 \langle PS; -\vec{p} | \mathcal{O}^{RL} | N; -\vec{k}, s \rangle \quad \Rightarrow \quad W_0^{LL} = W_0^{RR}, \\ W_0^{LR} = W_0^{RL}$$

- Relation on $[p,n] \rightarrow PS$ decay mode

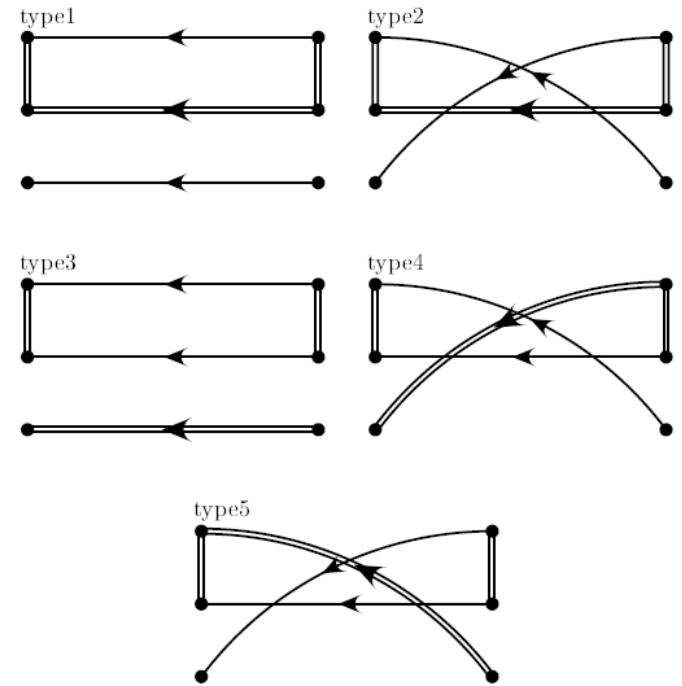
$$\begin{aligned} \langle \pi^0 | \mathcal{O}_{udu}^{L/R} | p \rangle &= \langle \pi^0 | \mathcal{O}_{dud}^{L/R} | n \rangle, \\ \langle \pi^+ | \mathcal{O}_{udd}^{L/R} | p \rangle &= -\langle \pi^- | \mathcal{O}_{duu}^{L/R} | n \rangle, \\ \langle K^0 | \mathcal{O}_{usu}^{L/R} | p \rangle &= -\langle K^+ | \mathcal{O}_{dsd}^{L/R} | n \rangle, \\ \langle K^+ | \mathcal{O}_{usd}^{L/R} | p \rangle &= -\langle K^0 | \mathcal{O}_{dsu}^{L/R} | n \rangle, \\ \langle K^+ | \mathcal{O}_{uds}^{L/R} | p \rangle &= -\langle K^0 | \mathcal{O}_{dus}^{L/R} | n \rangle, \\ \langle K^+ | \mathcal{O}_{dsu}^{L/R} | p \rangle &= -\langle K^0 | \mathcal{O}_{usd}^{L/R} | n \rangle, \\ \langle \eta | \mathcal{O}_{udu}^{L/R} | p \rangle &= -\langle \eta | \mathcal{O}_{dud}^{L/R} | n \rangle \end{aligned}$$

In the isospin limit $\langle \pi^+ | \mathcal{O}_{udu}^{L/R} | p \rangle = \sqrt{2} \langle \pi^0 | \mathcal{O}_{udu}^{L/R} | p \rangle$

\Rightarrow total number is **12**

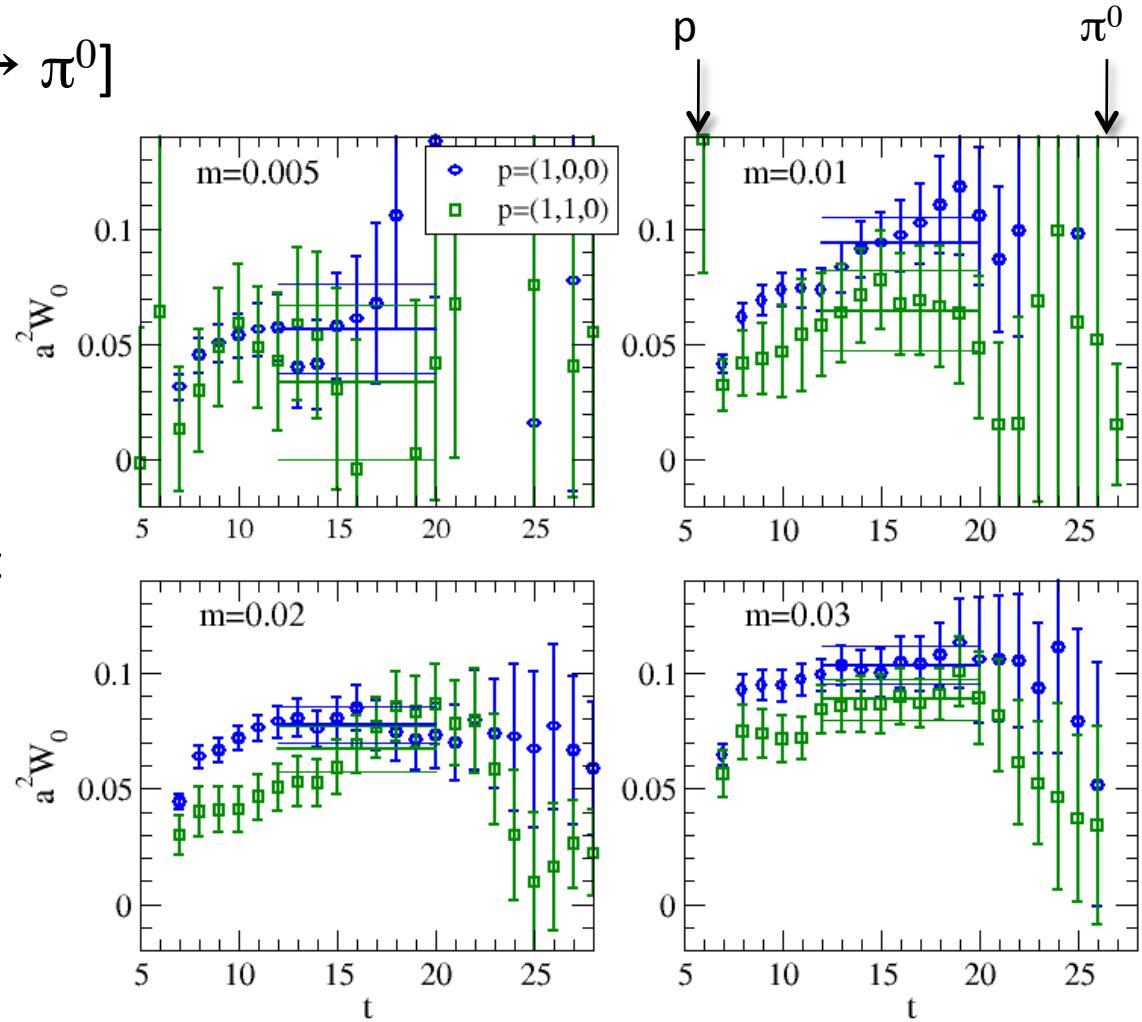
W_0 in direct method

$$\begin{aligned}
 \langle \pi^0 | \mathcal{O}_{udu} | p \rangle &= \frac{1}{\sqrt{2}} (-T^2 + T^3 + T^4 + T^5), \\
 \langle K^0 | \mathcal{O}_{usu} | p \rangle &= T^1 + T^2, \\
 \langle K^+ | \mathcal{O}_{usd} | p \rangle &= -T^2 + T^5, \\
 \langle K^+ | \mathcal{O}_{uds} | p \rangle &= T^3 + T^4, \\
 \langle K^+ | \mathcal{O}_{dsu} | p \rangle &= -T^4 - T^5, \\
 \langle \eta | \mathcal{O}_{udu} | p \rangle &= \frac{1}{\sqrt{6}} (2T^1 + T^2 + T^3 + T^4 + T^5),
 \end{aligned}$$



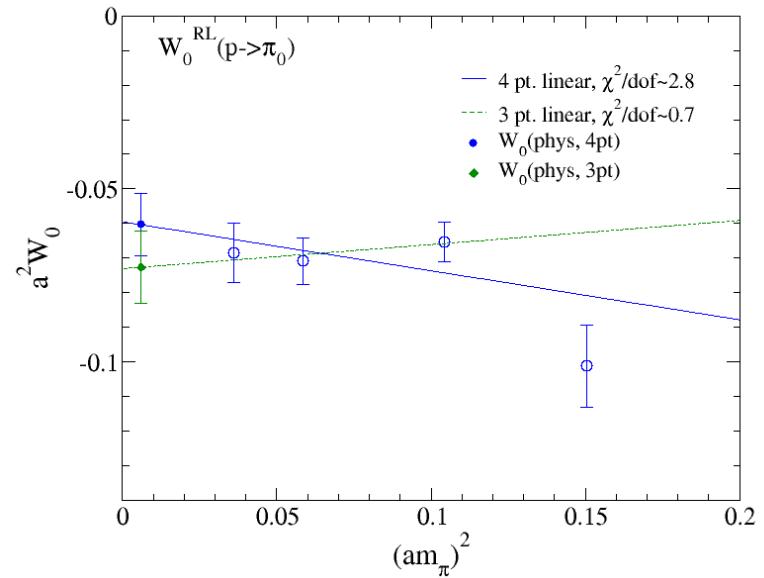
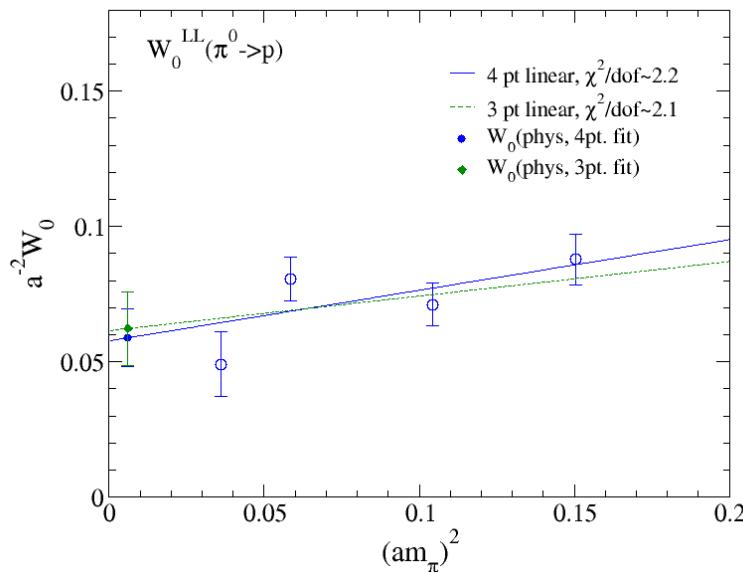
Results of W_0

- For $W_0^{\text{LL}}[p \rightarrow \pi^0]$
- Constant fit for plateau region
- Fit range [12,20]
- Large fluctuation at the lightest mass



Chiral extrapolation

- $W_0[p \rightarrow \pi^0]$ at $q^2 = 0$ point



- Linear extrapolation for 2 momentum
→ deterministic
- After that chiral extrapolation:

$$f(m_\pi^2, m_s) = W_0 + x_1 m_\pi^2 + x_2 (m_s - m_s^{\text{phy}})$$

RBC/UKQCD efforts

- Lattice size
 - 24^3 in spacial, 64 in temporal direction
 - $a^{-1} = 1.73 \text{ GeV}$,
 - 1.8^3 fm^3 size
 - Iwasaki-type gauge action + dynamical DW
- Domain-wall fermion
 - Chiral improved fermion by $L_s = 16$
 - Suppressed $O(a)$ error from χ breaking and unphysical operator mixing
 - Z factor from NPR Y. Aoki (2008)
 - $m=0.005, 0.01, 0.02, 0.03, m_s=0.03$
 - $m_\pi = 0.3 \text{ GeV} \sim 0.8 \text{ GeV}$

In the soft-pion limit

- Difference between W_0 [$p_\mu = 0$] and ChPT

